



Sub Topic : Spin Angular Momentum

1. What is spin angular momentum →

History : e^- 's spin (spin angular momentum) was discovered through Stern - Gerlach Experiment.

- Silver atoms (Ag) were used.
- Ag has 47 e^- 's. 46 of them form a spherically symmetric charge distribution and 47th e^- occupies 5s orbital.
- In ground state, angular momentum ($l=0$) must be zero for Ag as valence e^- is in 5s-orbital.
- In SG Experiment, non-homogeneous M.F is applied along z-axis.

According to Schrodinger model.

→ Beam of Ag atoms should go undeflected and we must get just one spot on screen. If Ag is in ground state (5s) $l=0$

→ If Ag is in excited state (5p) $l=1$. Beam will split into $2l+1 = 3$ beams and we get 3 spots on screen.



- Experimentally, beam splits into 2 components and we observe 2 spots on screen for $l=0$. Where no splitting was expected.
- **Conclusion:**
 - e^- possesses an intrinsic angular momentum which is called spin angular momentum.
 - It is analog to earth rotating around its own axis.
 - In real, e^- is point particle and it has no axis or rotation. e^- is born with this intrinsic spin.

Important:

- (1) Spin angular momentum does not depend upon its spatial DOF.
- (2) Spin is an intrinsic DOF.
- (3) Spin is purely QM concept which has no classical analog.
- (4) Unlike orbital angular momentum (L), spin can't be described by differential operator.



Magnetic Moment (μ_s) :-

- Magnetic moment of any body having angular momentum is

$$\vec{\mu}_s = -\frac{e \vec{S}}{2m} \rightarrow \text{classical prediction}$$

This would be true if $\bar{e}$ is assumed to be a charged sphere rotating around its own axis.

- In real, $\vec{\mu}_s = -2 \cdot \frac{e \vec{S}}{2m} = -g_s \frac{e \vec{S}}{2m}$

($g_s = 2$ Lande factor)

It is double the classical prediction.

External M.F $\rightarrow$

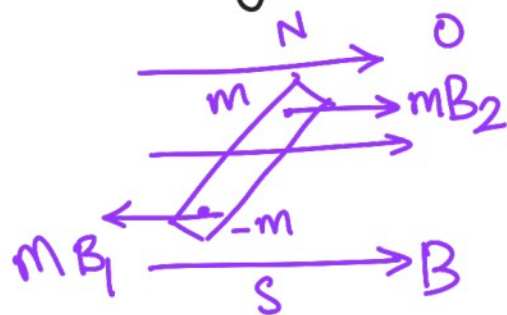
qE

- If $\bar{e}$ with magnetic moment μ_s is placed in external M.F (non-homogenous) it will experience force.

$$\vec{F} = \vec{\nabla} (\vec{\mu}_s \cdot \vec{B})$$

Why?

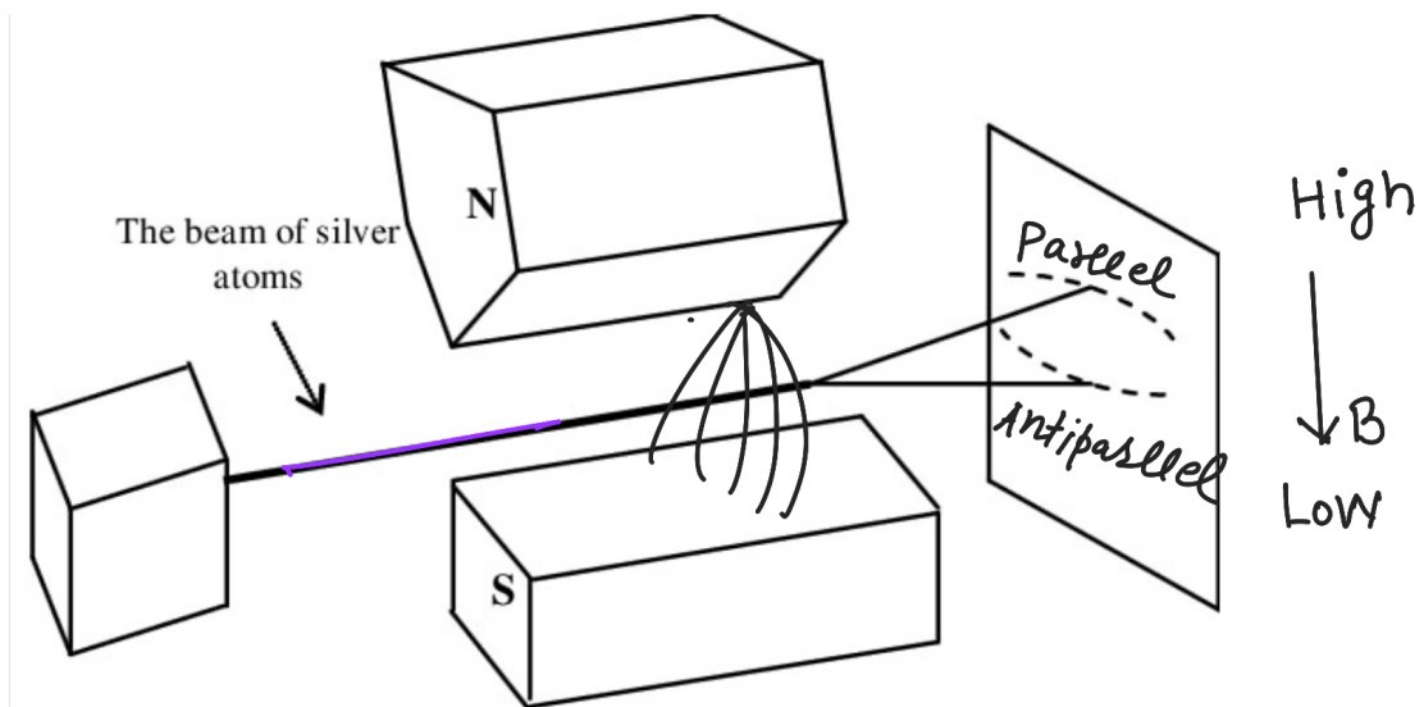
$$\vec{F} = -\vec{\nabla} u = \vec{\nabla} (\vec{\mu} \cdot \vec{B}) \quad (u = -\vec{\mu} \cdot \vec{B})$$





- Direction and magnitude of force depends upon relative orientation of field and dipole.
- If μ_s is in parallel to $M \cdot F$: $\vec{e}$ will move in direction in which field increase.
- If μ_s is antiparallel to $M \cdot F$: $\vec{e}$ will move in direction in which field decrease.

(I dare you to explain above 2 points)



Flat side magnet (low $M \cdot F$) field lines far
Curved " " (High $M \cdot F$) " " close



Orientation with M.F

- Spin of any particle is described by spin Q.N (S).
- $S = 1/2$ for $\bar{e}, p^+, n^0$
- Spin can have $2S + 1$ possible orientation w.r.t external M.F.
- Each orientation is represented by spin magnetic Q.N (m_s)
 $m_s = -S$ to S in steps of 1
- for $\bar{e}$, $S = 1/2$ $m_s = -1/2, 1/2$
- $\therefore \bar{e}$ beam splits into 2 beams.

$m_s = 1/2$	$m_s = -1/2$
Spin parallel to M.F	antiparallel.
- Spin $S = 3/2$ particle beam will split into $2S + 1 = 4$ beams ($m_s = -\frac{3}{2}, -\frac{1}{2}, \frac{1}{2}, \frac{3}{2}$)



2. Formalism of Spin $\rightarrow$

General Vector :

$$\vec{S} = \hat{S}_x \hat{i} + \hat{S}_y \hat{j} + \hat{S}_z \hat{k}$$

Commutator :

$$[\hat{S}_x, \hat{S}_y] = i\hbar \hat{S}_z$$

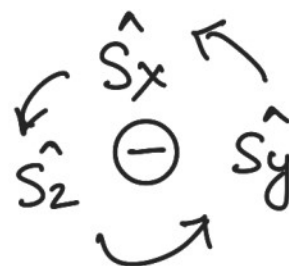
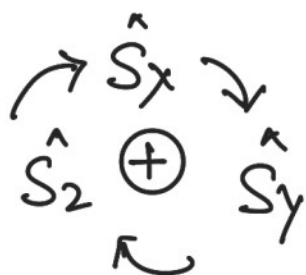
$$[\hat{S}_y, \hat{S}_x] = -i\hbar \hat{S}_z$$

$$[\hat{S}_y, \hat{S}_z] = i\hbar \hat{S}_x$$

$$[\hat{S}_z, \hat{S}_y] = -i\hbar \hat{S}_x$$

$$[\hat{S}_z, \hat{S}_x] = i\hbar \hat{S}_y$$

$$[\hat{S}_x, \hat{S}_z] = -i\hbar \hat{S}_y$$



$$[\hat{S}_x, \hat{S}_x] = [\hat{S}_y, \hat{S}_y] = [\hat{S}_z, \hat{S}_z] = 0$$

$\hat{S}_z$ and $\hat{S}^2$:

• $\hat{S}_z$ and $\hat{S}^2$ commute.

Significance

$$(1) [\hat{S}_z, \hat{S}^2] = 0$$

(2) Both operators have different



eigenvalues but same eigenvectors.

(3) Both can be measured simultaneously with 100% precision.

(4) Measurement of one operator don't disturb outcome of other.

(5) In general eigenvector of these operators are written as $|S, m_s\rangle$. It describes the state.

Eigenvalue equations

$$\hat{A} |S, m_s\rangle = \lambda |S, m_s\rangle$$

Diagram labels for the equation above:

- $\hat{A}$: operator
- $|S, m_s\rangle$ (left): eigenvector (State)
- λ : eigenvalue (measurement result)
- $|S, m_s\rangle$ (right): eigenvector (State)

$$\begin{aligned}\hat{S}^2 |S, m_s\rangle &= S(S+1)\hbar^2 |S, m_s\rangle \\ \hat{S}_z |S, m_s\rangle &= m_s \hbar |S, m_s\rangle\end{aligned}$$

* Eigenvalue for $\hat{S}^2$ is $S(S+1)\hbar^2$. If you measure $\hat{S}^2$ for system in state $|S, m_s\rangle$



You will get $S(S+1)\hbar^2$.

* Its eigenvalues of $\hat{S}_z$ is $m_s\hbar$.

Ladder operators :

$$\begin{cases} \hat{S}_+ = \hat{S}_x + i\hat{S}_y & \text{Increase } m_s \text{ by } +1 \\ \hat{S}_- = \hat{S}_x - i\hat{S}_y & \text{decrease } m_s \text{ by } -1 \end{cases}$$

$$\begin{cases} \hat{S}_+ |S, m_s\rangle = \sqrt{S(S+1) - m_s(m_s+1)}\hbar |S, m_s+1\rangle \\ \hat{S}_- |S, m_s\rangle = \sqrt{S(S+1) - m_s(m_s-1)}\hbar |S, m_s-1\rangle \end{cases}$$

Square operators :

• Average values : $\langle \hat{S}_x^2 \rangle = \langle \hat{S}_y^2 \rangle = ?$

$$\begin{cases} \langle \hat{S}_z^2 \rangle = m_s^2 \hbar^2 \\ \langle \hat{S}^2 \rangle = S(S+1) \hbar^2 \end{cases}$$

$$\langle \hat{S}^2 \rangle = \langle \hat{S}_x^2 \rangle + \langle \hat{S}_y^2 \rangle + \langle \hat{S}_z^2 \rangle$$

$$S(S+1)\hbar^2 = 2\langle \hat{S}_x^2 \rangle + m_s^2 \hbar^2$$



$$\langle \hat{S}_x^2 \rangle = \langle \hat{S}_y^2 \rangle = (S(S+1) - m_s^2) \frac{\hbar^2}{2}$$

Average of any operator :

- Average of operator $\hat{A}$ for system in state $|S, m_s\rangle$ is

$$\langle \hat{A} \rangle = \langle S, m_s | \hat{A} | S, m_s \rangle$$

- What does it mean?

System is in state $|S, m_s\rangle$. We measure physical quantity $\hat{A}$ for this state. We may get different value on each measurement. Average of physical quantity $\hat{A}$ is $\langle \hat{A} \rangle$.

Orthonormal states :

- Spin states form orthonormal basis

↓	↓
Orthogonal to each other	Normalized
$\langle S', m_{s'} S, m_s \rangle = 0 \quad \left(\begin{matrix} S' \neq S \\ m_{s'} \neq m_s \end{matrix} \right)$	$\langle S, m_s S, m_s \rangle = 1$



- They form complete basis.

$$\sum_{m_s=-s}^s |s, m_s\rangle \langle s, m_s| = I$$

Question: Jest 2022

Consider a spin-1 system whose $\hat{S}_z$ eigenstates are given by $|-1\rangle, |0\rangle, |1\rangle$ corresponding to the eigenvalues $-\hbar, 0, \hbar$. The normalized general state $|\psi\rangle$ of the system can be expressed as

$$|\psi\rangle = c_{-1} |-1\rangle + c_0 |0\rangle + c_{+1} |1\rangle,$$

and c_{-1}, c_0, c_{+1} are complex numbers. Subjected to the condition $\langle\psi|\hat{S}_z|\psi\rangle = 0$, which of the following statements is true?

(a) $|c_{+1}|^2 + 2|c_0|^2 = 1$

(c) $2|c_{-1}|^2 + |c_{+1}|^2 = 1$

(b) $|c_{-1}|^2 + 2|c_0|^2 = 1$

(d) $2|c_{-1}|^2 + |c_0|^2 = 1$

Solution:

$$S=1 \quad m_s = -1, 0, 1 \equiv |-1\rangle, |0\rangle, |1\rangle$$

- $\hat{S}_z |-1\rangle = -\hbar |-1\rangle$

$$\hat{S}_z |0\rangle = 0$$

$$\hat{S}_z |1\rangle = \hbar |1\rangle$$

- $|\psi\rangle = c_{-1} |-1\rangle + c_0 |0\rangle + c_{+1} |1\rangle$

- $\langle\psi|\hat{S}_z|\psi\rangle = ?$



$$\begin{aligned}
 \bullet \hat{S}_z |\psi\rangle &= \hat{S}_z (C_{-1} |-1\rangle + C_0 |0\rangle + C_1 |1\rangle) \\
 &= C_{-1} (-\hbar |-1\rangle) + 0 + C_1 (\hbar |1\rangle) \\
 &= \hbar (-C_{-1} |-1\rangle + C_1 |1\rangle)
 \end{aligned}$$

$$\bullet \langle \psi | \hat{S}_z | \psi \rangle = 0$$

$$\Rightarrow (-C_{-1}^* \langle -1| + C_0^* \langle 0| + C_1^* \langle 1|) (C_{-1} |-1\rangle + C_1 |1\rangle) \hbar = 0$$

$$\Rightarrow -C_{-1}^* C_{-1} + 0 + C_1^* C_1 = 0 \quad \left(\begin{array}{l} \text{only similar} \\ \text{terms give non} \\ \text{zero} \end{array} \right)$$

$$\Rightarrow |C_1|^2 - |C_{-1}|^2 = 0 \quad - (1)$$

From Normalization of Wavefunction

$$\langle \psi | \psi \rangle = 1$$

$$(C_{-1}^* \langle -1| + C_0^* \langle 0| + C_1^* \langle 1|) (C_{-1} |-1\rangle + C_0 |0\rangle + C_1 |1\rangle) = 1$$

$$|C_{-1}|^2 + |C_0|^2 + |C_1|^2 = 1 \quad - (2)$$



Subtract ① and ②

$$\boxed{2|c_{-1}|^2 + |c_0|^2 = 1} \quad \text{③}$$

Question: JEST 2023

Consider a spin-1/2 particle in the quantum state $|\psi(\beta, \alpha)\rangle = \cos\left(\frac{\beta}{2}\right)|\uparrow\rangle + \sin\left(\frac{\beta}{2}\right)e^{i\alpha}|\downarrow\rangle$ where $0 \leq \beta \leq \pi$ and $0 \leq \alpha \leq 2\pi$. For which values of (δ, γ) is the state $|\psi(\delta, \gamma)\rangle$ orthogonal to $|\psi(\beta, \alpha)\rangle$?

- (a) $(\pi + \beta, \pi - \alpha)$ (b) $(\pi - \beta, \pi - \alpha)$ (c) $(\pi + \beta, \pi + \alpha)$ (d) $(\pi - \beta, \pi + \alpha)$

Solution:

$$|\psi(\beta, \alpha)\rangle = \cos\left(\frac{\beta}{2}\right)|\uparrow\rangle + \sin\left(\frac{\beta}{2}\right)e^{i\alpha}|\downarrow\rangle$$

$$0 \leq \beta \leq \pi \quad \text{and} \quad 0 \leq \alpha \leq 2\pi$$

$$|\psi(\delta, \gamma)\rangle \text{ orthogonal to } |\psi(\beta, \alpha)\rangle$$

$$\langle \psi(\delta, \gamma) | \psi(\beta, \alpha) \rangle = 0$$

$$\langle \cos^*(\delta/2)\langle\uparrow| + \sin^*(\delta/2)e^{-i\gamma}\langle\downarrow| \rangle (\cos(\beta/2)|\uparrow\rangle + \sin(\beta/2)e^{i\alpha}|\downarrow\rangle) = 0$$

$$\cos^*(\delta/2) = \cos(\delta/2) \quad \left(\text{Prove it!}\right)$$

$$\sin^*(\delta/2) = \sin(\delta/2)$$

$$\cos\left(\frac{\delta}{2}\right)\cos\left(\frac{\beta}{2}\right) + \sin\left(\frac{\delta}{2}\right)\sin\left(\frac{\beta}{2}\right)e^{i(\alpha-\gamma)} = 0$$



Assume, $e^{i(\alpha-\gamma)} = -1$ ($= e^{i0}, e^{i(\pi)}, \dots$)

$$\alpha - \gamma = 0, \pi, \dots \Rightarrow \boxed{\gamma = \alpha - \pi}$$

$$\cos A \cos B = \frac{1}{2} (\cos(A+B) + \cos(A-B))$$

$$\sin A \sin B = \frac{1}{2} (\cos(A-B) - \cos(A+B))$$

$$\therefore \frac{1}{2} \left(\cos\left(\frac{\delta+\beta}{2}\right) + \cos\left(\frac{\delta-\beta}{2}\right) - \cos\left(\frac{\delta+\beta}{2}\right) + \cos\left(\frac{\delta-\beta}{2}\right) \right) = 0$$

$$\therefore \cos\left(\frac{\delta-\beta}{2}\right) = 0$$

$$\frac{\delta-\beta}{2} = \frac{\pi}{2} \Rightarrow \delta - \beta = \pi$$

$$\boxed{\delta = \beta + \pi}$$

$$\boxed{\gamma = \pi - \alpha}$$

if we do $\langle \psi(\beta, \alpha) | \psi(\delta, \gamma) \rangle = 0$

(Option a)



3. Matrix Representation $\rightarrow$

$$S = 1/2 \text{ particle}$$

$$\bullet S = 1/2 \quad m_s = \pm 1/2$$

$$\bullet \text{Possible states } |S, m_s\rangle = |1/2, 1/2\rangle, |1/2, -1/2\rangle$$

$\hookrightarrow |\uparrow\rangle$ $\hookrightarrow |\downarrow\rangle$

$\hat{S}^2$ - Matrix

$$\hat{S}^2 = \begin{pmatrix} \langle \uparrow | \hat{S}^2 | \uparrow \rangle & \langle \uparrow | \hat{S}^2 | \downarrow \rangle \\ \langle \downarrow | \hat{S}^2 | \uparrow \rangle & \langle \downarrow | \hat{S}^2 | \downarrow \rangle \end{pmatrix}$$

$$\hat{S}^2 |\uparrow\rangle = \hbar^2 S(S+1) |\uparrow\rangle = \frac{3}{4} \hbar^2 |\uparrow\rangle$$

$$\hat{S}^2 |\downarrow\rangle = \frac{3}{4} \hbar^2 |\downarrow\rangle$$

$$\hat{S}^2 = \begin{pmatrix} \frac{3\hbar^2}{4} \langle \uparrow | \uparrow \rangle & \frac{3\hbar^2}{4} \langle \uparrow | \downarrow \rangle \\ \frac{3\hbar^2}{4} \langle \downarrow | \uparrow \rangle & \frac{3\hbar^2}{4} \langle \downarrow | \downarrow \rangle \end{pmatrix}$$

$$\hat{S}^2 = \frac{3\hbar^2}{4} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$



$\hat{S}_z$ - Matrix :

$$\hat{S}_z = \begin{pmatrix} \langle \uparrow | \hat{S}_z | \uparrow \rangle & \langle \uparrow | \hat{S}_z | \downarrow \rangle \\ \langle \downarrow | \hat{S}_z | \uparrow \rangle & \langle \downarrow | \hat{S}_z | \downarrow \rangle \end{pmatrix}$$

$$\hat{S}_z |\uparrow\rangle = \frac{\hbar}{2} |\uparrow\rangle \quad \text{and} \quad \hat{S}_z |\downarrow\rangle = -\frac{\hbar}{2} |\downarrow\rangle$$

$$\hat{S}_z = \begin{pmatrix} \hbar/2 & 0 \\ 0 & -\hbar/2 \end{pmatrix}$$

$$\boxed{\hat{S}_z = \frac{\hbar}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}}$$

$\hat{S}_+$ Matrix :

$$\hat{S}_+ = \begin{pmatrix} \langle \uparrow | \hat{S}_+ | \uparrow \rangle & \langle \uparrow | \hat{S}_+ | \downarrow \rangle \\ \langle \downarrow | \hat{S}_+ | \uparrow \rangle & \langle \downarrow | \hat{S}_+ | \downarrow \rangle \end{pmatrix}$$

$$\hat{S}_+ |\uparrow\rangle = \hat{S}_+ \left| \frac{1}{2}, \frac{1}{2} \right\rangle = \sqrt{S(S+1) - m_s(m_s+1)} \hbar = 0$$

$$\hat{S}_+ |\downarrow\rangle = \hat{S}_+ \left| \frac{1}{2}, -\frac{1}{2} \right\rangle = \sqrt{\frac{3}{4} + \frac{1}{4}} \hbar \left| \frac{1}{2}, \frac{1}{2} \right\rangle = \hbar |\uparrow\rangle$$

$$\hat{S}_+ = \begin{pmatrix} 0 & \hbar \langle \uparrow | \uparrow \rangle \\ 0 & \hbar \langle \downarrow | \uparrow \rangle \end{pmatrix} = \begin{pmatrix} 0 & \hbar \\ 0 & 0 \end{pmatrix}$$



$$\hat{S}_+ = \hbar \begin{pmatrix} 0 & 1 \\ 0 & 1 \end{pmatrix}$$

$\hat{S}_-$ Matrix :

$$\hat{S}_- = \begin{pmatrix} \langle \uparrow | \hat{S}_- | \uparrow \rangle & \langle \uparrow | \hat{S}_- | \downarrow \rangle \\ \langle \downarrow | \hat{S}_- | \uparrow \rangle & \langle \downarrow | \hat{S}_- | \downarrow \rangle \end{pmatrix}$$

$$\begin{aligned} \hat{S}_- | \uparrow \rangle &= \hat{S}_- | \tfrac{1}{2}, \tfrac{1}{2} \rangle = \sqrt{S(S+1) - m_s(m_s-1)} \hbar | \tfrac{1}{2}, -\tfrac{1}{2} \rangle \\ &= \sqrt{\tfrac{3}{4} + \tfrac{1}{4}} \hbar | \tfrac{1}{2}, -\tfrac{1}{2} \rangle \end{aligned}$$

$$\hat{S}_- | \uparrow \rangle = \hbar | \downarrow \rangle$$

$$\hat{S}_- | \downarrow \rangle = 0$$

$$\hat{S}_- = \begin{pmatrix} \langle \uparrow | \downarrow \rangle \hbar & 0 \\ \langle \downarrow | \downarrow \rangle \hbar & 0 \end{pmatrix} = \begin{pmatrix} 0 & 0 \\ \hbar & 0 \end{pmatrix}$$

$$\hat{S}_- = \hbar \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix}$$



$\hat{S}_x$ Matrix :

$$\hat{S}_x = \frac{1}{2} (\hat{S}_+ + \hat{S}_-) = \frac{\hbar}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$\hat{S}_y = \frac{i}{2} (\hat{S}_- - \hat{S}_+) = \frac{\hbar}{2} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$

$$\hat{S}_z = \frac{\hbar}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

Pauli - Matrices :

$$\sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \quad \sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \quad \sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

Properties

(1) $\sigma_x^2 = \sigma_y^2 = \sigma_z^2 = I$ (involuntary matrices)

(2)

Commutators

$$[\hat{A}, \hat{B}] = \hat{A}\hat{B} - \hat{B}\hat{A}$$

$$[\sigma_x, \sigma_y] = 2i\sigma_z$$

$$[\sigma_y, \sigma_z] = 2i\sigma_x$$

$$[\sigma_z, \sigma_x] = 2i\sigma_y$$

Anticommutators

$$\{\hat{A}, \hat{B}\} = \hat{A}\hat{B} + \hat{B}\hat{A}$$

$$\{\sigma_x, \sigma_y\} = 0$$

$$\{\sigma_y, \sigma_z\} = 0$$

$$\{\sigma_z, \sigma_x\} = 0$$



(3) Determinant = -1

(4) Trace = 0

(5) Eigenvalues = ± 1

(6) Eigenvectors

$$\psi_{x+} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix} \quad \psi_{x-} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -1 \end{pmatrix}$$

$$\psi_{y+} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ i \end{pmatrix} \quad \psi_{y-} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -i \end{pmatrix}$$

$$\psi_{z+} = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \quad \psi_{z-} = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

Spin Matrices

$$S_x = \frac{\hbar}{2} \sigma_x$$

$$S_y = \frac{\hbar}{2} \sigma_y$$

$$S_z = \frac{\hbar}{2} \sigma_z$$

$$\vec{S} = \frac{\hbar}{2} \vec{\sigma}$$

Formulae to help problem solving :

$$\textcircled{1} (\vec{\sigma} \cdot \vec{A})(\vec{\sigma} \cdot \vec{B}) = (\vec{A} \cdot \vec{B}) \hat{I} + i \vec{\sigma} \cdot (\vec{A} \times \vec{B})$$

$$\textcircled{2} \sigma_x \sigma_y \sigma_z = i \hat{I}$$

$$\textcircled{3} e^{i\alpha \sigma_j} = \hat{I} \cos \alpha + i \sigma_j \sin \alpha \quad (j=x, y, z)$$



commutators with other quantities $\rightarrow$

- Spin does not depend on spatial DOF, so spin commute with all spatial operators

$$[\hat{S}_j, \hat{L}_k] = 0 \quad [\hat{S}_j, \hat{R}_k] = 0 \quad [\hat{S}_j, \hat{P}_k] = 0$$

$$\text{eg } [S_x, L_z] = 0 \quad [S_y, z] = 0 \quad [S_x, P_y] = 0$$

Spin operators in terms of pauli matrices

For $S = 1/2$

$$S_x = \frac{\hbar}{2} \sigma_x$$

$$S_y = \frac{\hbar}{2} \sigma_y$$

$$S_z = \frac{\hbar}{2} \sigma_z$$

$$\pm \hbar/2$$

eigenvalues
 $\pm \hbar/2$

$$\pm \hbar/2$$

eigenvectors

$$\frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix}, \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -1 \end{pmatrix}$$

$$\frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ i \end{pmatrix}, \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -i \end{pmatrix}$$

$$\begin{pmatrix} 1 \\ 0 \end{pmatrix}, \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$



Eigenvectors of $\hat{S}_z$ as basis vectors :

• $\hat{S}_z$ has 2 eigenvectors $|\uparrow\rangle_z = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$

$$|\downarrow\rangle_z = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

• In general state of system can be written as :

$$\Psi = \alpha |\uparrow\rangle_z + \beta |\downarrow\rangle_z$$

Probability to get $|\uparrow\rangle_z$ (Spin up) = $\frac{\alpha^2}{\alpha^2 + \beta^2}$

" " " $|\downarrow\rangle_z$ (Spin down) = $\frac{\beta^2}{\alpha^2 + \beta^2}$

• How can we write our state in terms of eigenvectors of $\hat{S}_x$?

$$|\uparrow\rangle_x = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix}$$

$$|\downarrow\rangle_x = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -1 \end{pmatrix}$$



we need to write $|\uparrow\rangle_z$, $|\downarrow\rangle_z$ in terms of $|\uparrow\rangle_x$ and $|\downarrow\rangle_x$

$$\text{let, } |\uparrow\rangle_z = a |\uparrow\rangle_x + b |\downarrow\rangle_x$$

$$a = \langle \uparrow | \uparrow \rangle_z = \frac{1}{\sqrt{2}} (1 \ 1) \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \frac{1}{\sqrt{2}}$$

$$b = \langle \downarrow | \uparrow \rangle_z = \frac{1}{\sqrt{2}} (1 \ -1) \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \frac{1}{\sqrt{2}}$$

$$\therefore \boxed{|\uparrow\rangle_z = \frac{1}{\sqrt{2}} (|\uparrow\rangle_x + |\downarrow\rangle_x)}$$

$$\text{liky, let } |\downarrow\rangle_z = a |\uparrow\rangle_x + b |\downarrow\rangle_x$$

$$a = \langle \uparrow | \downarrow \rangle_z = \frac{1}{\sqrt{2}} (1 \ 1) \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \frac{1}{\sqrt{2}}$$

$$b = \langle \downarrow | \downarrow \rangle_z = \frac{1}{\sqrt{2}} (1 \ -1) \begin{pmatrix} 0 \\ 1 \end{pmatrix} = -\frac{1}{\sqrt{2}}$$

$$\therefore \boxed{|\downarrow\rangle_z = \frac{1}{\sqrt{2}} (|\uparrow\rangle_x - |\downarrow\rangle_x)}$$



$$\therefore \psi = \alpha |\uparrow\rangle_z + \beta |\downarrow\rangle_z$$

$$\psi = \frac{\alpha}{\sqrt{2}} (|\uparrow\rangle_x + |\downarrow\rangle_x) + \frac{\beta}{\sqrt{2}} (|\uparrow\rangle_x - |\downarrow\rangle_x)$$

$$\psi = \left(\frac{\alpha+\beta}{\sqrt{2}}\right) |\uparrow\rangle_x + \left(\frac{\alpha-\beta}{\sqrt{2}}\right) |\downarrow\rangle_x$$

$$\text{Probability to get } |\uparrow\rangle_x = \frac{\left(\frac{\alpha+\beta}{\sqrt{2}}\right)^2}{\left(\frac{\alpha+\beta}{\sqrt{2}}\right)^2 + \left(\frac{\alpha-\beta}{\sqrt{2}}\right)^2}$$

$$\text{Probability to get } |\downarrow\rangle_x = \frac{\left(\frac{\alpha-\beta}{\sqrt{2}}\right)^2}{\left(\frac{\alpha+\beta}{\sqrt{2}}\right)^2 + \left(\frac{\alpha-\beta}{\sqrt{2}}\right)^2}$$

- How can we write our state in terms of eigenvectors of $\hat{S}_y$.

$$|\uparrow\rangle_y = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ i \end{pmatrix}$$

$$|\downarrow\rangle_y = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -i \end{pmatrix}$$



$$\text{let, } |\uparrow\rangle_z = a |\uparrow\rangle_y + b |\downarrow\rangle_y$$

$$a = \langle \uparrow | \uparrow \rangle_z = \frac{1}{\sqrt{2}} (1 - i) \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \frac{1}{\sqrt{2}}$$

$$b = \langle \downarrow | \uparrow \rangle_z = \frac{1}{\sqrt{2}} (1 \ i) \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \frac{1}{\sqrt{2}}$$

$$\therefore \boxed{|\uparrow\rangle_z = \frac{1}{\sqrt{2}} (|\uparrow\rangle_y + |\downarrow\rangle_y)}$$

$$\text{let, } |\downarrow\rangle_z = a |\uparrow\rangle_y + b |\downarrow\rangle_y$$

$$a = \langle \uparrow | \downarrow \rangle_z = \frac{1}{\sqrt{2}} (1 - i) \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \frac{-i}{\sqrt{2}}$$

$$b = \langle \downarrow | \downarrow \rangle_z = \frac{1}{\sqrt{2}} (1 \ i) \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \frac{i}{\sqrt{2}}$$

$$\boxed{|\downarrow\rangle_z = \frac{i}{\sqrt{2}} (-|\uparrow\rangle_y + |\downarrow\rangle_y)}$$

$$\therefore \psi = \alpha |\uparrow\rangle_z + \beta |\downarrow\rangle_z$$

$$\psi = \frac{\alpha}{\sqrt{2}} (|\uparrow\rangle_y + |\downarrow\rangle_y) + \frac{\beta i}{\sqrt{2}} (-|\uparrow\rangle_y + |\downarrow\rangle_y)$$



$$\psi = \frac{(\alpha - i\beta)}{\sqrt{2}} |\uparrow\rangle_y + \left(\frac{\alpha + i\beta}{\sqrt{2}}\right) |\downarrow\rangle_y$$

Question: Spin $1/2$ particle is in state

$$\chi = \frac{1}{\sqrt{6}} \begin{pmatrix} 1+i \\ 2 \end{pmatrix}$$

a) If you measure S_z , what is probability that you will get $\pm \hbar/2$.

$$\chi = \frac{1+i}{\sqrt{6}} \begin{pmatrix} 1 \\ 0 \end{pmatrix} + \frac{2}{\sqrt{6}} \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

$$\chi = \alpha |\uparrow\rangle_z + \beta |\downarrow\rangle_z \quad \left(\begin{array}{l} \alpha = \frac{1+i}{\sqrt{6}} \\ \beta = \frac{2}{\sqrt{6}} \end{array} \right)$$

$$\alpha^2 + \beta^2 = \left| \frac{1+i}{\sqrt{6}} \right|^2 + \left| \frac{2}{\sqrt{6}} \right|^2 = \frac{2}{6} + \frac{4}{6} = 1$$

$$P(=+\hbar/2) = \left| \frac{1+i}{\sqrt{6}} \right|^2 = \frac{2}{6} = \frac{1}{3}$$

$$P(=-\hbar/2) = \left| \frac{2}{\sqrt{6}} \right|^2 = \frac{4}{6} = \frac{2}{3}$$



b) If you measure $\hat{S}_x$. What is the probability of getting $\pm \hbar/2$.

$$P(=\hbar/2) = \left| \frac{\alpha + \beta}{\sqrt{2}} \right|^2 = \left| \frac{3+i}{\sqrt{2}} \right|^2 = \frac{10}{2} = \frac{5}{1}$$

$$P(=-\hbar/2) = \left| \frac{\alpha - \beta}{\sqrt{2}} \right|^2 = \left| \frac{-1+i}{\sqrt{2}} \right|^2 = \frac{2}{2} = \frac{1}{1}$$

c) What is expectation value of $\hat{S}_x$:

$$\langle \hat{S}_x \rangle = \langle \psi | \hat{S}_x | \psi \rangle$$

$$= \begin{pmatrix} \frac{1-i}{\sqrt{2}} & \frac{2}{\sqrt{2}} \end{pmatrix} \begin{pmatrix} 0 & \hbar/2 \\ \hbar/2 & 0 \end{pmatrix} \begin{pmatrix} (1+i)/\sqrt{2} \\ 2/\sqrt{2} \end{pmatrix}$$

$$= \frac{\hbar}{2} \begin{pmatrix} \frac{1-i}{\sqrt{2}} & \frac{2}{\sqrt{2}} \end{pmatrix} \begin{pmatrix} 2/\sqrt{2} \\ (1+i)/\sqrt{2} \end{pmatrix}$$

$$= \frac{\hbar}{2} \left(\frac{2-2i+2+2i}{2} \right) = \frac{\hbar}{2} \frac{4}{2} = \frac{\hbar}{1}$$



Spin Operator in any random direction

• Find eigenvalues and eigenvectors of spin operator $\vec{S}$ of an electron in the direction of unit vector $\vec{n}$ ($\vec{n}$ is arbitrary).

• Operator : $\vec{n} \cdot \vec{S}$

• Eigenvalue equation :

$$\vec{n} \cdot \vec{S} |\psi\rangle = \frac{\hbar}{2} \lambda |\psi\rangle$$

• $\vec{n}$ in spherical coordinate system

$$\vec{n} = \sin\theta \cos\phi \hat{i} + \sin\theta \sin\phi \hat{j} + \cos\theta \hat{k}$$

$$\vec{S} = S_x \hat{i} + S_y \hat{j} + S_z \hat{k}$$

$$\vec{n} \cdot \vec{S} = (S_x \sin\theta \cos\phi) + (S_y \sin\theta \sin\phi) + (S_z \cos\theta)$$



$$\vec{n} \cdot \vec{S} = \frac{\hbar}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \sin\theta \cos\phi + \frac{\hbar}{2} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \sin\theta \sin\phi + \frac{\hbar}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \cos\theta$$

$$= \frac{\hbar}{2} \begin{pmatrix} \cos\theta & \sin\theta (\cos\phi - i\sin\phi) \\ \sin\theta (\cos\phi + i\sin\phi) & -\cos\theta \end{pmatrix}$$

$$= \frac{\hbar}{2} \begin{pmatrix} \cos\theta & \sin\theta e^{-i\phi} \\ \sin\theta e^{i\phi} & -\cos\theta \end{pmatrix}$$

eigenvalues

$$\begin{vmatrix} \frac{\hbar}{2} \cos\theta - \lambda & \frac{\hbar}{2} \sin\theta e^{-i\phi} \\ \frac{\hbar}{2} \sin\theta e^{i\phi} & -\frac{\hbar}{2} \cos\theta - \lambda \end{vmatrix} = 0$$

$$-\frac{\hbar^2}{4} (\cos^2\theta - \lambda^2) - \frac{\hbar^2}{4} \sin^2\theta = 0$$

$$\frac{\lambda^2 \hbar^2}{4} - \frac{\hbar^2}{4} = 0$$



$$\lambda^2 = 1$$

$$\lambda = \pm 1$$

eigenvectors :

$$|\lambda\rangle_+ = \begin{pmatrix} \cos\theta/2 \\ e^{i\phi} \sin\theta/2 \end{pmatrix}$$

$$|\lambda\rangle_- = \begin{pmatrix} -\sin\theta/2 \\ e^{i\phi} \cos\theta/2 \end{pmatrix}$$

Example : Test 2023

Consider the operator $S \cdot \hat{n}$ with eigenkets $|\pm\rangle_{\hat{n}}$ and eigenvalues $\pm \frac{\hbar}{2}$ where $\hat{n}$ is a unit vector and S is the spin operator. A partially polarized beam of spin- $\frac{1}{2}$ particles contains a 25-75 mixture of two pure ensembles, one with $|+\rangle_{\hat{z}}$ and the other with $|+\rangle_{\hat{x}}$ respectively. What is the ensemble average of

$$\frac{S \cdot \hat{x}}{\hbar}?$$

- (a) $\frac{1}{3}$
- (b) $\frac{1}{8}$
- (c) $\frac{1}{4}$
- (d) $\frac{3}{16}$



$$\frac{\hat{S} \cdot \hat{x}}{\hbar} \text{ is simply } \frac{S_x}{\hbar}$$

our system is in state

$$\psi = \frac{1}{2} |\uparrow\rangle_z + \frac{\sqrt{3}}{2} |\uparrow\rangle_x$$

$$\psi = \frac{1}{2} \begin{pmatrix} 1 \\ 0 \end{pmatrix} + \frac{\sqrt{3}}{2\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix} = \frac{1}{2\sqrt{2}} \begin{pmatrix} \sqrt{2} + \sqrt{3} \\ \sqrt{3} \end{pmatrix}$$

average of $S_x/\hbar$

$$\frac{\langle \psi | S_x | \psi \rangle}{\hbar} = \frac{\hbar}{2\hbar} \frac{1}{8} (\sqrt{2} + \sqrt{3} \quad \sqrt{3}) \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} \sqrt{2} + \sqrt{3} \\ \sqrt{3} \end{pmatrix}$$

$$= \frac{1}{16} (\sqrt{2} + \sqrt{3} \quad \sqrt{3}) \begin{pmatrix} \sqrt{3} \\ \sqrt{2} + \sqrt{3} \end{pmatrix}$$

$$= \frac{1}{16} (2\sqrt{3}(\sqrt{2} + \sqrt{3}))$$

$$= \frac{\sqrt{3}(\sqrt{2} + \sqrt{3})}{8}$$

ans not matching



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